Biophysical Model Building

- **Step 1:** Come up with a hypothesis about how a system works
 - How many binding sites?
 - Is there cooperativity?
- Step 2: Translate the qualitative hypotheses into an observable mathematical form with parameters
 - Example parameters: K, tau, N
 - Parameters may not be known
- **Step 3:** Design an experiment that that can produce observables from step 2; perform the experiment
 - Optimize the parameters to make the fit look as good as possible
- Step 4: Assess the fit Is the agreement convincing?

Example: Single Site Binding

- Situation: You are studying a novel DNA-binding protein
- Hypothesis: $P + D \leftrightarrow PD$
 - Parameters: K
 - Implicitly, we assume that N = 1, no cooperativity
- Experiment: Collect a binding curve (dialysis)
 - Optimize K for the best fit
- Assess: How good is our fit? Use statistics!

Example: Single Site Binding

• Si What about the Thermodynamics?

• Free energies:

$$\Delta \bar{G}^0 = -RT \ln K$$

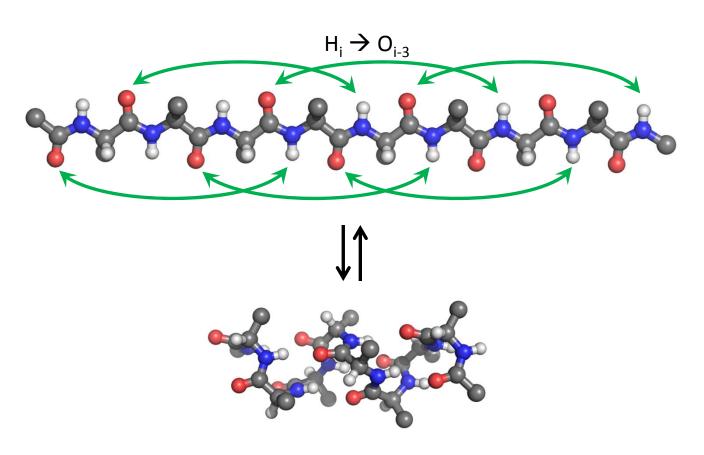
Enthalpies: Van't Hoff

$$\ln \frac{K_2}{K_1} = -\frac{\Delta \overline{H}^0}{R} \left(\frac{1}{T_2} - \frac{1}{T_1} \right)$$

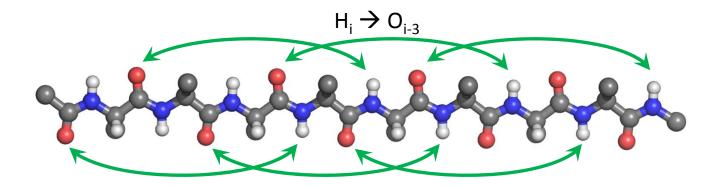
• Once the appropriate thermodynamic variables are known, one can predict K at any T, P, etc.

Helix-Coil Theory

• Question: How does an α -helix fold?



Helix-Coil Theory: Considerations



- Entropic cost: 3 pairs of φ, ψ torsions must be "fixed" before 1 hydrogen bond is formed
- Energetic benefit: forming H-bond is favorable once torsions are "fixed"
- End effects: No H-bonds for the end residues

- Assumption 1: Each residue can exist in one of two conformational states: h or c
 - h (helix) is one conformation, c (coil) represents many energetically equivalent conformations
 - We can "enumerate" confirmations with a series of h's and c's
 - For 8 residues:

hhhhcccc ccchhccc hchchchc hhhhhhh

- Assumption 2: For any sequence we consider, assume it's flanked by an infinite number of "coil" residues
 - This allows us to ignore end effects

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ccchhccc → ...ccccchhcccccc...

hhhhhhhh → ...ccchhhhhhhhhccc...

hhhhcccc → ...ccchhhhhccccccc...
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- Assumption 3: Assume some conformations are not observed
 - hch, hcch: helical kinks, but no energetic benefit

ccchhccc	\rightarrow	possible
hhhhhhhh	\rightarrow	possible
hchhhccc	\rightarrow	not possible
hhhcchhh	\rightarrow	not possible
hhccchhh	\rightarrow	possible

- Assumption 4a: Individual residues are in equilibrium
 - If a new helical residue tries to form at the end of an existing helix:

 $hhhccccc \rightarrow hhhhcccc K_{eq} = s$

- Assumption 4b: Individual residues are in equilibrium
 - If a new helical residue tries to form "from scratch":

$$ccccccc \rightarrow ccchcccc K_{eq} = \sigma s$$

- σ represents entropic cost of forming a helical φ, ψ values ($\sigma \ll 1$)

- Assumption 5: Let the "unfolded" (i.e. all-coil) state be our reference state
 - Statistical weight of ...ccccccc... = 1

Make a table (for N = 3):

State	# Helical Residues	Weight
<i>ccc</i>	0	$\frac{[ccc]}{[ccc]} = 1$
hcc	1	$\frac{[hcc]}{[ccc]} = ?$
chc	1	$\frac{[chc]}{[ccc]} = ?$
cch	1	$\frac{[cch]}{[ccc]} = ?$
hhc	2	$\frac{[hhc]}{[ccc]} = ?$
hch	2	$\frac{[hch]}{[ccc]} = ?$
chh	2	$\frac{[chh]}{[ccc]} = ?$
hhh	3	$\frac{[hhh]}{[ccc]} = ?$

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cch	1	$\frac{[cch]}{[ccc]} = \sigma s$
hhc	2	$\frac{[hhc]}{[ccc]} = ?$
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chh	2	$\frac{[chh]}{[ccc]} = \sigma s^2$
hhh	3	$\frac{[hhh]}{[ccc]} = ?$

Disallowed state!

Make a table (for N = 3):

State	# Helical Residues	Weight
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chc	1	$\frac{[chc]}{[ccc]} = \sigma s$
cch	1	$\frac{[cch]}{[ccc]} = \sigma s$
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hch	2	$\frac{[hch]}{[ccc]} = 0$
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hhh	3	$\frac{[hhh]}{[ccc]} = \sigma s^3$

Helix-Coil Theory: Experiment

Partition function:

$$Z = \sum w_i$$

$$Z = 1 + 3\sigma s + 2\sigma s^2 + \sigma s^3$$

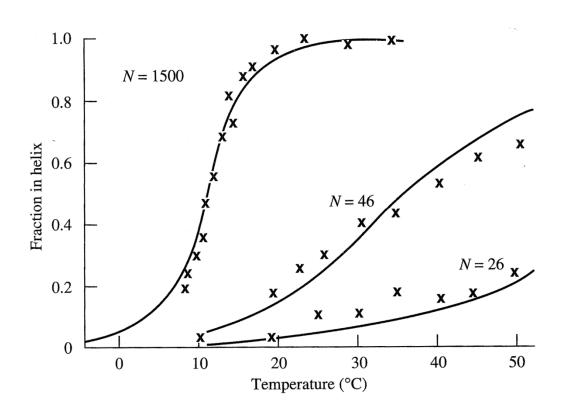
 Weighted average # helical residues:

$$\langle \text{helix} \rangle = \frac{\sum n_i w_i}{Z}$$

<helix> is measurable!

State	# Helical Residues	Weight
ccc	0	$\frac{[ccc]}{[ccc]} = 1$
hcc	1	$\frac{[hcc]}{[ccc]} = \sigma s$
chc	1	$\frac{[chc]}{[ccc]} = \sigma s$
cch	1	$\frac{[cch]}{[ccc]} = \sigma s$
hhc	2	$\frac{[hhc]}{[ccc]} = \sigma s^2$
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hhh	3	$\frac{[hhh]}{[ccc]} = \sigma s^3$

Helix-Coil Theory: Experiment



◄ FIGURE 11.7

Temperature dependence of the fraction of monomer units in the helix form for poly-γ-benzyl-L-glutamate (in a 7:3 mixture of dichloracetic acid and 1,2-dichloroethane as the solvent) with polymer lengths of 26, 46, and 1500 monomer units. In this solvent the helix is favored at high temperatures. (Data from B. H. Zimm and J. K. Bragg, 1959, J. Chem. Phys. 31:526; B. H. Zimm, P. Doty, and K. Iso, 1959, Proc. Natl. Acad. Sci. USA 45:1601.)

Helix-Coil Theory: Takeaways

- Typical parameters
 - $-s \approx 1$ (sometimes favorable, sometimes unfavorable depending on the solvent \rightarrow hydrogen bond strength?
 - $-\sigma \ll 1$ (typically on the order 0.001) \rightarrow entropy is difficult to overcome
- Cooperativity comes from small σ relative to s
- It's hard to break a helix in the middle
 - Helix typically "frays" at the ends

Helix-Coil Theory: DNA?

- DNA is more complicated than protein
 - Different base pairs → different s values
 - Two strands involved $\rightarrow \sigma$ is concentration dependent
 - Individual strands can "shift" relative to one another

Lots of people have studied this (see book)

Random Walk in 1D

- Molecule diffuses l (average distance) before colliding with another molecule
 - This is the "step size" per hop
- For one step, probability (weight) of traveling forward is p, probability of traveling backward is q = (1 p)

Random Walk for 3 Steps

Partition function:

$$Z = \sum w_i Z = p^3 + 3p^2q + 3pq^2 + q^3$$

Avg. # forward steps:

$$\langle \mathbf{m} \rangle = \frac{\sum m_i w_i}{Z}$$

Mean displacement:

$$\langle \mathbf{d} \rangle = [\langle m \rangle - (N - \langle m \rangle)]l$$

 $\langle \mathbf{d} \rangle = [2\langle m \rangle - N]l$

• When
$$p = \frac{1}{2}$$
, $\langle d \rangle = 0$
- Why?

Steps	# Forward Steps	Weight
fff	3	p^3
bff	2	p^2q
fbf	2	p^2q
ffb	2	p^2q
bbf	1	pq^2
bfb	1	pq^2
fbb	1	pq^2
bbb	0	q^3

Random Walk

 Random walk in 1-D can be solved in general for N steps (see p. 631-33)

• Main point: When $p = \frac{1}{2}$, the *mean-squared* displacement is:

$$\langle d^2 \rangle = N l^2$$

 This can be used to model the end-to-end distance of an unfolded chain with no self-avoidance (N = # of links)

What Have We Learned?

- Many biological systems are made up of discrete states
 - Bound → Free; Native → Unfolded; Helix → Coil
- If we can determine relative concentrations (weights), we can sum the weights to get a partition function
- Once we know weights and the partition function, we can calculate observable values
 - Mole fractions: $X_i = \frac{w_i}{Z}$
 - Weighted averages: $\langle x \rangle = \frac{\sum x_i w_i}{Z}$
 - Degree of binding: $\bar{\nu} = \frac{\sum n_i w_i}{Z}$

What Don't We Know?

• Can we calculate weights directly from energy differences (i.e. $\Delta \bar{G}$)?

 Can we relate entropy to the number of possible states?

What is the molecular significance of the partition function?